地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」 2021.03.04

Adjoint-based exact Hessian computation

Shin-ichi Ito^{1,2}

Today's talk is based on **BIT Numerical Mathematics (2021)**



01/30

Joint work with

Takeru Matsuda³ and Yuto Miyatake⁴

- **1. Earthquake Research Institute, The University of Tokyo**
- 2. Graduate School of Information Science and Technology, The University of Tokyo
- **3. RIKEN Center for Brain Science**
- 4. Cybermedia Center, Osaka University

We developed an adjoint-based method

that enables exact Hessian computations up to a given floating point limit.

Error in a Hessian matrix Column Column Row Row **Our method** Naive method

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Outline

Introduction Adjoint-based gradient and Hessian computations Sanz-Serna's idea and exact Hessian computation Numerical verifications

02/30

Introduction Adjoint-based gradient and Hessian computations

03/30

Adjoint method

We consider an autonomous ordinary differential equation(ODE):

$$x_0 = \theta$$
 \rightarrow $\frac{d}{dt} x_t = f(x_t)$ \rightarrow $x_T(\theta)$
Forward model

and a minimization problem of a cost function defined by the final state t = T

$$\min_{\theta} C\left(x_T\left(\theta\right)\right)$$

via an gradient method using a gradient $\
abla_{ heta} C$

Adjoint method calculates $abla_{ heta} C$ exactly by solving



$$\lambda_0 = \nabla_\theta C \qquad \longleftarrow \qquad -\frac{d}{dt} \lambda_t = (\nabla_x f)^\top \lambda_t \qquad \longleftarrow \qquad \lambda_T = \nabla_{x_T} C$$
Adjoint model

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Second-order adjoint method



* These model are derived from the first-order perturbation of the forward and adjoint models.

06/30

© Shin-ichi Ito

SOA provides an exact Hessian-vector product for arbitrary vectors.

SOA is available for optimization methods, uncertainty quantification, ...

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Application

- Fast optimization
 - Newton method $\theta \leftarrow \theta + y$ where $H_{\theta}y = -\nabla_{\theta}C$ **V: Descent vector**

 $C(\theta)$

Nonlinear conjugate gradient method





「固体地球現象の理解と予測に向けたデータ同化法の開発」 地震研特定研究 (B)

Question

Adjoint and SOA methods give exact gradient $\nabla_{\theta}C$ and Hessian-vec. prod. $H_{\theta}\gamma$ if all of the models are solved analytically.

In practice, all of the models are solved numerically,

using time integral schemes, e.g., Runge–Kutta methods.

→ inexact due to discretization errors, machine errors



$$\frac{d}{dt}x_t = f\left(x_t\right)$$

Tangent linear (TL) model

$$\frac{d}{dt}\delta_t = (\nabla_x f)\,\delta_t$$

Adjoint model

 $-\frac{d}{dt}\lambda_t = (\nabla_x f)^\top \lambda_t$

SOA model

$$-\frac{d}{dt}\xi_t = (\nabla_x f)^\top \xi_t + \{(\nabla_x \nabla_x f) \,\delta_t\}^\top \,\lambda_t$$

For a given forward scheme, are there "optimal schemes" for other models?

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Question

Adjoint and SOA methods give exact gradient $\nabla_{\theta}C$ and Hessian-vec. prod. $H_{\theta}\gamma$ if all of the models are solved analytically.

In practice, all of the models are solved numerically,

using time integral schemes, e.g., Runge–Kutta methods.

→ inexact due to discretization errors, machine errors

| Forward model | | |
|--|--|--|
| $\frac{d}{dt}x_t = f\left(x_t\right)$ | | |
| Adjoint model | | |
| $-\frac{d}{dt}\lambda_t = (\nabla_x f)^\top \lambda_t$ | | |
| Sanz-Serna (2016) — | | |

Tangent linear (TL) model

$$\frac{d}{dt}\delta_t = (\nabla_x f)\,\delta_t$$

SOA model

$$-\frac{d}{dt}\xi_t = (\nabla_x f)^\top \xi_t + \{(\nabla_x \nabla_x f) \delta_t\}^\top \lambda_t$$

For a given forward scheme, are there "optimal schemes" for other models? \rightarrow Yes. There exists an optimal scheme for the adjoint model that can avoid any discretization errors. 09/30

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Question

Adjoint and SOA methods give exact gradient $abla_{\theta}C$ and Hessian-vec. prod. $ext{H}_{\theta}\gamma$ if all of the models are solved analytically.

In practice, all of the models are solved numerically,

using time integral schemes, e.g., Runge–Kutta methods.

→ inexact due to discretization errors, machine errors



For a given forward scheme, are there "optimal schemes" for other models?

→ Yes. There exists a set of optimal schemes for all of these models that can avoid any discretization errors.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Importance of exact Hessian computation

***** Numerical details will be explained later.

Error in a Hessian matrix



地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Importance of exact Hessian computation



When using an incorrect scheme,

- Optimization may not converge to correct solution.
- Uncertainty quantification may fail.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Sanz-Serna's idea and exact Hessian computation

Invariant

Why can the analytical solution of the adjoint model provide the exact gradient?

A product of TL and adjoint variables is time-invariant

TL model $\frac{d}{dt}\delta_t = (\nabla_x f)\delta_t$

- **Adjoint model** -
$$\frac{d}{dt}\lambda_t = (\nabla_x f)^\top \lambda_t$$

14/30

If we choose $\lambda_t = \nabla_{x_t} C(x_t(\theta))$, we obtain

$$\delta_t^\top \nabla_{x_t} C(x_t) = ((\nabla_\theta x_t) \,\delta_0)^\top \,\nabla_{x_t} C(x_t)$$
$$= \delta_0^\top \left[(\nabla_\theta x_t)^\top \,\nabla_{x_t} C(x_t) \right] = \delta_0^\top \nabla_\theta C(x_t)$$

$$\lambda_0 = \nabla_\theta C(x_t(\theta))$$

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

$$\therefore x_t (\theta + \epsilon \delta_0) = x_t(\theta) + \epsilon \delta_t + O(\epsilon^2)$$
$$\flat \delta_t = (\nabla_\theta x_t) \delta_0$$

Runge–Kutta family



* (...) : order

15/30

| Classical 4-stage RK (4) | • Explicit Euler (1) | • Heun (2) |
|---|-----------------------------|---|
| $x_{n+1} = x_n + \frac{h}{6} \left(k_{n,1} + 2k_{n,2} + 2k_{n,3} + k_{n,4} \right),$ | $x_{n+1} = x_n + hk_{n,1},$ | $x_{n+1} = x_n + \frac{h}{2} \left(k_{n,1} + k_{n,2} \right),$ |
| $k_{n,i} = f(X_{n,i})$ $(i = 1,, 4),$ | $k_{n,1}=f(X_{n,1}),$ | $k_{n,i} = f(X_{n,i})$ $(i = 1,, 2),$ |
| $X_{n,1} = x_n$ $X_{n,3} = x_n + \frac{n}{2}k_{n,2}$ | $X_{n,1} = x_n$ | $X_{n,1} = x_n$ |
| $X_{n,2} = x_n + \frac{h}{2}k_{n,1} \qquad X_{n,4} = x_n + hk_{n,3}$ | | $X_{n,2} = x_n + hk_{n,1}$ |

 Adaptive stepsize RK: Bogacki-Shampine(2-3), Runge-Kutta-Fehlberg(4-5), Dormand-Prince(4-5), ...

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Sanz-Serna's idea

Solve forward models via a s-stage Runge–Kutta (RK) method

and solve adjoint model via another s-stage RK method

Sanz-Serna (2016) provided a relation between (a, b) and (A, B)

so that

$$\delta_{n+1}^{\top}\lambda_{n+1} = \delta_n^{\top}\lambda_n$$
 exactly.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Sanz-Serna's idea

$$\delta_{n+1}^{\mathsf{T}}\lambda_{n+1} = \delta_n^{\mathsf{T}}\lambda_n$$

$$b_i = B_i$$
 (*i* = 1,..., *s*),
 $b_i A_{ij} + B_j a_{ji} = B_i b_j$ (*i*, *j* = 1,..., *s*).

$\delta_{n+1} = \delta_n + h \sum_{i=1}^{n} b_i d_{n,i},$ Proof $d_{n,i} = \nabla_x f(X_{n,i}) D_{n,i} \quad (i = 1, \ldots, s),$ $S(\delta_t, \lambda_t) = \delta_t^{\top} \lambda_t$ $D_{n,i} = \delta_n + h \sum_{i=1}^{n} a_{ij} d_{n,j}$ (i = 1, ..., s). $S(\delta_{n+1},\lambda_{n+1}) - S(\delta_n,\lambda_n)$ $= h \sum_{i=1}^{5} B_i S(\delta_n, l_{n,i}) + h \sum_{i=1}^{5} b_i S(d_{n,i}, \lambda_n) + h^2 \sum_{i=1}^{5} B_i b_j S(d_{n,i}, l_{n,j})$ $\lambda_{n+1} = \lambda_n + h \sum B_i l_{n,i},$ $l_{n,i} = -\nabla_x f(X_{n,i})^\top \Lambda_{n,i} \quad (i = 1, \dots, s),$ $= h \sum_{i=1}^{n} B_i S(D_{n,i}, l_{n,i}) - h^2 \sum_{i=1}^{n} B_i a_{ij} S(d_{n,j}, l_{n,i})$ $+h\sum_{i=1}^{s}b_{i}S(d_{n,i},\Lambda_{n,i})-h^{2}\sum_{i,i=1}^{s}b_{i}A_{ij}S(d_{n,i},l_{n,j})+h^{2}\sum_{i,i=1}^{s}B_{i}b_{j}S(d_{n,i},l_{n,j}) \qquad \Lambda_{n,i}=\lambda_{n}+h\sum_{j=1}^{s}A_{ij}l_{n,j} \quad (i=1,\ldots,s).$ $=h\sum_{i=1}^{s} (b_{i} - B_{i}) S\left(D_{n,i}, \nabla_{x} f(X_{n,i})^{\top} \Lambda_{n,i}\right) + h^{2} \sum_{i,j=1}^{s} (B_{i}b_{j} - b_{i}A_{ij} - B_{j}a_{ji}) S(d_{n,i}, l_{n,j})$

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Exact gradient computation

Forward models solved by a s-stage Runge–Kutta (RK) method

and adjoint model solved by another s-stage RK method

$$-\frac{d}{dt}\lambda_t = (\nabla_x f)^\top \lambda_t \quad \bullet \quad I_{n,i} = -\nabla_x f(X_{n,i})^\top \Lambda_{n,i} \quad (i = 1, \dots, s),$$
$$\Lambda_{n,i} = \lambda_n + h \sum_{j=1}^s A_{ij} l_{n,j} \quad (i = 1, \dots, s).$$

where

$$b_i = B_i$$
 (*i* = 1,..., *s*),
 $b_i A_{ij} + B_j a_{ji} = B_i b_j$ (*i*, *j* = 1,..., *s*).

provide exact gradient as $\lambda_0 = \nabla_{\theta} C$

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Adjoint consistency in time



In general,

G₁ and G₂ are different.

Neither G_1 nor G_2 gives an exact gradient (G_0)

because $\delta_{n+1}^{\top}\lambda_{n+1} \neq \delta_n^{\top}\lambda_n$ in both cases.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Adjoint consistency in time



「固体地球現象の理解と予測に向けたデータ同化法の開発」 地震研特定研究 (B)

Towards exact Hessian computation

Exact gradient can be computed by using Sanz-Serna's scheme.

Exact Hessian ???

Forward model $\frac{d}{dt}x_t = f(x_t)$ Adjoint model $-\frac{d}{dt}\lambda_t = (\nabla_x f)^\top \lambda_t$ Sanz-Serna (2016)

Tangent linear (TL) model

$$\frac{d}{dt}\delta_t = (\nabla_x f)\,\delta_t$$

SOA model

$$-\frac{d}{dt}\xi_t = (\nabla_x f)^\top \xi_t + \{(\nabla_x \nabla_x f) \,\delta_t\}^\top \,\lambda_t$$

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Towards exact Hessian computation

Our idea: Equivalence between SOA model and adjoint model

$$\frac{d}{dt}x_{t} = f(x_{t}) \qquad \frac{d}{dt}\delta_{t} = (\nabla_{x}f)\delta_{t}$$
$$-\frac{d}{dt}\lambda_{t} = (\nabla_{x}f)^{\top}\lambda_{t} \qquad -\frac{d}{dt}\xi_{t} = (\nabla_{x}f)^{\top}\xi_{t} + \{(\nabla_{x}\nabla_{x}f)\delta_{t}\}^{\top}\lambda_{t}$$
$$Define augmented vectors: \quad q_{t} = \begin{bmatrix} x_{t} \\ \delta_{t} \end{bmatrix} \quad \text{and} \quad p_{t} = \begin{bmatrix} \lambda_{t} \\ \xi_{t} \end{bmatrix}$$
$$\frac{d}{dt}q_{t} = \begin{bmatrix} f(x_{t}) \\ (\nabla_{x}f)\delta_{t} \end{bmatrix} = G(q_{t})$$

$$-\frac{d}{dt}p_t = \begin{bmatrix} (\nabla_x f)^\top \lambda_t \\ (\nabla_x f)^\top \xi_t + \{(\nabla_x \nabla_x f) \delta_t\}^\top \lambda_t \end{bmatrix} = (\nabla_q G)^\top p_t$$

Set of adjoint and SOA models constitutes a large adjoint system

→ Many techniques for adjoint method can be used for SOA method. e.g., Sanz-Serna's idea, adjoint code generator

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Exact gradient and Hessian computations

The augmented models solved by <u>Sanz-Serna (2016</u>), i.e., the augmented forward model solved by a s-stage RK method

and the augmented adjoint model solved by another s-stage RK method

$$p_{n+1} = p_n + h \sum_{i=1}^{s} B_i l_{n,i},$$

$$l_{n,i} = -\nabla_q G(Q_{n,i})^\top P_{n,i} \quad (i = 1, ..., s),$$

$$P_{n,i} = p_n + h \sum_{j=1}^{s} A_{ij} l_{n,j} \quad (i = 1, ..., s).$$
where
$$b_i = B_i \quad (i = 1, ..., s),$$

$$b_i A_{ij} + B_j a_{ji} = B_i b_j \quad (i, j = 1, ..., s).$$

provide exact gradient and Hessian computations.

Note that our method itself is not limited to RK.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Numerical verifications - Inhomogeneous wave equation

© Shin-ichi Ito 23/30

Inhomogeneous wave equation

Validate our method through numerical test of data assimilation based on a one-dimensional inhomogeneous wave equation

$$\frac{\partial^2 U(z,t)}{\partial t^2} = \frac{\partial}{\partial z} \left[E(z) \frac{\partial}{\partial z} U(z,t) \right]$$

(discretized in space by finite volume method)

Time integral scheme:

Heun method (2-stage RK)

Assume initial states of U and \dot{U} are known but E is not.

Cost function:

$$C(E) = \sum_{t_n \in T^{\text{obs}}} ||U_n(E) - U_n^{\text{obs}}||_2^2$$
, where



$$T^{\text{obs}} = \{0, 0.2, 0.4, \dots, 1.8, 2.0\}$$

© Shin-ichi Ito

Discretization via finite volume method (staggered lattice)

$$\begin{aligned} \frac{\mathrm{d}}{\mathrm{d}t}U_i &= V_i \quad (i = 1, \dots, d) \\ \frac{\mathrm{d}}{\mathrm{d}t}V_i &= \frac{1}{\Delta z^2} \begin{cases} E_{\frac{3}{2}}\left(U_2 - U_1\right) - E_{d+\frac{1}{2}}\left(U_1 - U_d\right) & (i = 1) \\ E_{i+\frac{1}{2}}\left(U_{i+1} - U_i\right) - E_{i-\frac{1}{2}}\left(U_i - U_{i-1}\right) & (i = 2, \dots, d-1) \\ E_{d+\frac{1}{2}}\left(U_1 - U_d\right) - E_{d-\frac{1}{2}}\left(U_d - U_{d-1}\right) & (i = d) \end{cases} \\ \\ \frac{\mathrm{d}}{\mathrm{d}t}E_{i+\frac{1}{2}} &= 0 \quad (i = 1, \dots, d), \end{aligned}$$

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Time integral schemes



地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

Symmetry of Hessian matrix

 $||H_{\theta} - H_{\theta}^{\top}||_{\infty}$: Degree of asymmetry, which is zero if H_{θ} is symmetric.



Our method reproduces the symmetry of Hessian without any discretization errors even when using a large step size.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Convergence to optimum solution



Our method contributes to rapid convergence.

地震研特定研究 (B) 「固体地球現象の理解と予測に向けたデータ同化法の開発」

© Shin-ichi Ito

Large-scale uncertainty quantification

<u>Slow-slip motion model</u>





A, B, and L are spatially dependent Ito et al., in prep.

Time integrator: Runge-Kutta-Fehlberg method



* We proposed a method to compute an exact Hessian-vector product.

* The fact that the second-order adjoint system can be reformulated to a part of a large adjoint system is the key point to obtain the exact Hessian-vector product based on the Sanz-Serna scheme.

* Our method is capable of being applied to various types of ODEs.

* Seismic imaging, aerodynamic design, structural materials, ...

* Neural ODE

* Thank you for your attention



k{0/5



30/30